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# Molecular Complexes Exhibiting Polarization Bonding. XI. The Crystal and Molecular Structure of the 7,7,8,8-Tetracyanoquinodimethane-Anthracene Complex

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The 1:1 complex formed between 7,7,8,8-tetracyanoquinodimethane and anthracene crystallizes as deep green needles, space group C2/m, with two molecules of complex in the unit cell of dimensions  $a = 11.476 \pm 0.011$ ,  $b = 12.947 \pm 0.013$ ,  $c = 7.004 \pm 0.007$  Å,  $\beta = 105.4 \pm 0.2^{\circ}$ . The structure was solved by Patterson methods and refined by least-squares methods using 469 independent reflexions measured on a linear diffractometer. The final residual, R, was 0.111. The component molecules are stacked alternately in infinite columns which are parallel to the c axis with the mean plane of the molecules perpendicular to the stacking axis and a mean perpendicular separation of 3.50 Å. The central six-membered rings of the two types of molecule are orientated at 30° to each other.

# Introduction

The 1:1 complex formed between 7,7,8,8-tetracyanoquinodimethane (TCNQ) and anthracene was first described by Melby, Harder, Hertler, Mahler, Benson & Mochel (1962). The investigation of the crystal and molecular structure of the complex was carried out to provide comparison with the structures of complexes of aromatic molecules with tetracyanoethylene (TCNE) as exemplified by the naphthalene–TCNE complex described in Part IX (Williams & Wallwork, 1967).

### Experimental

Deep-green needles of the complex were deposited after hot equimolar solutions of the components in tetrahydrofuran were mixed and allowed to cool very slowly. Microanalyses for carbon, nitrogen and hydrogen confirmed the 1:1 ratio of the component molecules in the complex. In plane polarized light the crystals exhibited strong pleochroism. The colours ranged from deep green to a straw colour as the plane of polarization of the light was rotated and maximum absorption occurred when the electric vector of the polarized light was vibrating parallel to the needle axis.

Oscillation and Weissenberg X-ray photographs were taken, using Cu  $K\alpha$  radiation ( $\lambda = 1.542$  Å), with the crystal rotating about the *b* and the *c* (needle) axes. The cell dimensions, determined from the Weissenberg photographs, were refined with the use of a linear diffractometer (Arndt & Phillips, 1961), and intensity data for the h0l, h1l, ..., h11l and the hk0, hk1, ..., hk6 reciprocal lattice levels were then collected with Mo K $\alpha$ radiation. A total of 469 independent reflexions with  $\theta < 32.5^{\circ}$  were observed out of a possible total of 1800 accessible reflexions. Intensity measurements were made with a scintillation counter and a pulse height selector. The moving-crystal stationary-counter technique was employed. The crystal used was approximately cylindrical with a radius of 0.2 mm. No absorption corrections were applied to the data.

# Crystal data

 $C_8H_4(CN)_4$ .  $C_{14}H_{10}$ , M = 382.43. Monoclinic,  $a = 11.476 \pm 0.011$ ,  $b = 12.947 \pm 0.013$ ,  $c = 7.004 \pm 0.007$  Å,  $\beta = 105.4^{\circ} \pm 0.2^{\circ}$ , U = 1003 Å<sup>3</sup>,  $D_m = 1.25$  g.cm<sup>-3</sup>, Z = 2,  $D_c = 1.26$  g.cm<sup>-3</sup>, F(000) = 396. Mo  $K\alpha$  ( $\lambda = 0.7107$  Å),  $\mu = 0.84$  cm<sup>-1</sup>. Absent spectra, hkl when h+k odd. Space group C2 (no.5), Cm (no.8) or C2/m (no.12). Negative piezoelectric and pyroelectric tests indicate C2/m and this is confirmed by refinement of the structure.

# Determination and refinement of the structure

A trial structure was deduced from the hk0 and h0lPatterson projections. The structure was partially refined through structure-factor, electron-density projection calculations to the stage where the reliability index,

$$R = \Sigma ||F_o| - |F_c|| / \Sigma |F_o|,$$

was 0.21 and 0.16 for the hk0 and h0l zones respectively. Structure factors were then calculated for the hkl reflexions followed by a three-dimensional Fourier electron-density summation, from which improved atomic coordinates were obtained.

Refinement was continued by the block-diagonal least-squares method with a program written by Dr R.D.Diamand. The weighting scheme employed was  $\sqrt[]{w=1}$  if  $|F_o| < F^*$  otherwise  $\sqrt[]{w=F^*/|F_o|}$ , and the scattering factors of Hoerni & Ibers (1954) for carbon and nitrogen were used. The reflexions too weak to be observed were omitted from the analysis. Hydrogen atom parameters were not included. Four cycles of structure-factor least-squares using individual isotropic temperature factors improved R to 0.165. Refinement

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was then continued allowing the temperature factors of the atoms to vary individually and anisotropically. After four cycles of least-squares the refinement converged with R=0.111 for observed reflexions. The comparison of observed and final calculated structure factors is shown in Table 6.

The final atomic coordinates and their estimated standard deviations (derived from the least-squares normal equations block-diagonal matrix) are summarized in Table 1. The anisotropic temperature factors are given in Table 2.

Table 1. Final positional parameters (upper line) and their estimated standard deviations (lower line)

	$10^{4}x/a$	104 <i>y/b</i>	$10^{4}z/c$
C(1)	0000	1088	0000
	—	9	
C(2)	1077	0551	0475
	7	6	12
C(3)	2214	1127	0938
	10	12	16
C(4)	3228	0499	1382
	9	13	15
C(5)	1271	0000	5509
	9	—	14
C(6)	0617	0975	5241
	6	5	11
C(7)	2506	0000	6026
	9		15
C(8)	3198	0929	6325
	6	6	12
Ν	3764	1657	6574
	6	6	12

Table 2. Final atomic thermal parameters ( $\times 10^4$ )

 $B_{ij}$  are coefficients in the temperature factor expression exp  $[-(B_{11}h^2 + B_{22}k^2 + B_{33}l^2 + B_{12}hk + B_{13}hl + B_{23}kl)].$ 

	$B_{11}$	B <sub>22</sub>	B <sub>33</sub>	$B_{12}$	$B_{13}$	B <sub>23</sub>
C(1)	198	87	367	000	205	000
C(2)	145	122	367	- 55	140	- 37
C(3)	230	297	443	-214	235	-116
C(4)	217	432	433	-152	156	-65
C(5)	116	69	268	000	129	000
C(6)	132	71	306	13	124	9
C(7)	122	72	318	000	83	000
C(8)	114	91	442	-1	123	28
N	149	100	634	- 54	120	-3

Sections of the three-dimensional electron-density distribution and of the difference electron-density distribution were then calculated. These syntheses showed no unexpected features, confirmed the marked anisotropy of the anthracene molecule and clearly indicated the positions of the hydrogen atoms of the TCNQ molecule.

Least-squares planes through each of the molecules were calculated by the modification (Blow, 1960) of the method of Schomaker, Waser, Marsh & Bergman (1959). All of the atoms, other than hydrogen, in each molecule were used to define the planes and each atom used was given unit weight. The pertinent parameters of the planes, and the deviations of the atoms from them, are summarized in Table 3.

Table 3. Molecular least-squares planes

	TCNQ	Anthracene
Perp	endicular dist	ance in Å from the origin
	3.502	0.000
Direction	cosines of nor	mal relative to orthogonal axes
а	0.2533	0.2612
Ь	0.0000	0.0000
с*	-0.9674	- 0.9653
mandiaula	distance in d	of atoms from the mean play

Perpendicular distance in Å of atoms from the mean planes

C(5)	0.013	C(1)	0.000
C(6)	0.010	C(2)	-0.010
C(7)	0.009	C(3)	0.006
C(8)	0.001	C(4)	-0.001
N(1)	-0.008		

The anisotropic thermal parameters of the atoms were used in an analysis of the rigid-body translational and librational motion (Cruickshank, 1956) of each molecule. The results of the analyses are given in Table 4 for the TCNQ molecule and in Table 5 for the anthracene molecule. The r.m.s. discrepancies between the observed and calculated  $U_{ij}$  values are 0.004 and 0.017 Å<sup>2</sup> for the TCNQ and anthracene molecules respectively. The method of Busing & Levy (1964) was used to correct the bond lengths of the TCNQ molecule for libration. The corrected bond lengths are illustrated in Fig. l(a) together with the uncorrected bond lengths and angles and their e.s.d.'s. The bond lengths and angles and their e.s.d.'s for the anthracene molecule are illustrated in Fig. 1(b). No attempt was made to correct these for librational motion because it was so extreme as to cause atoms C(4) and C(4') to be unresolved in one electron-density peak.

#### Description and discussion of the structure

The structure consists of infinite columns of alternate TCNO and anthracene molecules stacked in a planeto-plane manner in columns parallel to the c axis of the unit cell with the mean molecular planes perpendicular to this axis. Each cell contains two columns of molecules and the structure, projected along the caxis, is illustrated in Fig.2. The mode of molecular overlap is such that the long axes of each molecule are parallel, and their centres are superimposed. The central hexagonal ring systems of the two types of molecule are orientated at 30° to each other, thereby avoiding direct overlap of the majority of the atoms. This relative orientation is not that for the optimum overlap of the  $\pi$ -orbital systems of the molecules but the structure is such that the aromatic system of the TCNQ molecule is overlapped by the entire aromatic system of the anthracene molecule. There is no suggestion of specific interaction involving C = C bonds as found in the structure of the naphthalene-TCNE complex (Williams & Wallwork, 1967).

The most surprising aspect of the structure is the relatively large intermolecular spacing (3.50 Å) of the



Fig. 1. Bond lengths (Å) and bond angles, with (in parentheses) their standard deviations, for (a) TCNQ, (b) anthracene. Values for TCNQ in square brackets are bond lengths corrected for libration.

# Table 4. Rigid-body thermal parameters for the TCNQ molecule (For the molecule centred at $0, 0, \frac{1}{2}$ ).

$$\tau^{\dagger} = \begin{pmatrix} 746 & -1 & 729 \\ 547 & 2 \\ 536 \end{pmatrix} \times 10^{-4} \text{ Å}^2 \qquad \omega^{\dagger} = \begin{pmatrix} 32 \cdot 5 & 0 \cdot 1 & 8 \cdot 1 \\ 9 \cdot 3 & -0 \cdot 7 \\ 8 \cdot 1 \end{pmatrix} (^{\circ})^2$$

Direction cosines ( $\times 10^4$ ) for principal axes relative to orthogonal axes:

Eigenvalue τ 0.0769 Å <sup>2</sup> 0.0547	a 9543 139	<i>b</i> -7 -9988	<i>c</i> * 2987 469		Eigenvalue $\omega$ 34.8 (°) <sup>2</sup> 9.6	9 	2 596 450	-32 -9888	c* 2814 1425
0.0213	2984	489	9593	(\$)2 - 10	5.8	-2	778	1494	9490
			$O_{ij}$	$(A)^{2} \times \Pi$	<b>1</b>				
		$U_{11}$	$U_{22}$	$U_{33}$	$U_{12}$	$U_{13}$	$U_{23}$		
	C(5)	79	56	71	-2	3	0		
		1	4	0	6	2	2		
	C(6)	75	59	59	0	6	0		
		-7	1	3	0	2	0		
	C(7)	76	68	75	0	2	1		
		2	-7	-2	0	-6	-1		
	C(8)	80	77	102	-8	-5	2		
	- (- )	9	0	0	6	1	9		
	N(1)	89	86	147	- 16	-16	2		
	- (1)	8	-1	-1	-4	0	$-\bar{3}$		
		•	<u> </u>	-		•	5		

† Referred to the directions of the orthogonal axes  $a, b, c^*$ .

<sup>‡</sup> The numbers below the  $U_{ij}$  are the differences (× 10<sup>3</sup>) of  $U_{ij}$  derived from  $B_{ij}$  and those calculated from rigid-body parameters. The r.m.s. difference is 0.004 (Å)<sup>2</sup>.

molecules within one stack, presumably due to the unfavourable relative orientation of the components. The four shortest approaches between atoms of adjacent molecules of the same column (C(1)–C(6) 3.57, C(2)– C(6) 3.58, C(2)–C(5) 3.58, C(4)–C(8) 3.58 Å) are, if anything, greater than the expected van der Waals distances. This implies that any charge-transfer interaction between the molecules is weak. This does not conflict with the existence of a low-energy, charge-transfer band in the spectrum of the complex, which causes its deep green colour. The fairly low ionization potential



Fig. 2. The molecular arrangement, viewed along the c axis.

of anthracene and the high electron affinity of TCNQ inevitably produce an excited state of fairly low energy, but this will not stabilize the ground state of the complex by admixture of the two wave functions if the symmetries are inappropriate for overlap.

Both molecules are planar within experimental error and the molecular planes are inclined to each other at the small angle of  $\frac{1}{2}^{\circ}$ . The structure projected along the *b* axis is illustrated in Fig. 3.

The libration-corrected dimensions of the TCNQ molecule are in good agreement with the dimensions of the uncomplexed molecule (Long, Sparks & Trueblood, 1965). The r.m.s. discrepancy between the two determinations is 0.007 Å, much less than one e.s.d. No bond length differs by more than  $1\frac{1}{2}$  e.s.d. The maximum difference in the bond angles is also  $1\frac{1}{2}$  e.s.d., the r.m.s. difference being one degree.

The dimensions of the anthracene molecule [apart from the bonds and angles involving the atom C(4)] are in reasonable agreement with the weighted mean of the bond lengths determined by various authors and discussed by Cruickshank & Sparks (1960). The dimensions involving the atom C(4) are invalid for the reasons discussed above.

The thermal parameters,  $U_{ij}$  of the TCNQ molecule appear to be adequately explained on the basis of rigidbody librations and vibrations. The principal axes of both libration and of vibration correspond to within a few degrees of the natural molecular axes and there is appreciable libration only about the long axis in the plane of the molecule. The thermal parameters,  $U_{ij}$ , of the anthracene molecule are reasonably well explained on the basis of rigid-body motions. The dis-

Table 5.	Rigid-body	parameters f	for the	anthracene	molecule
	(For th	e molecule cer	ntred at	: 0,0,0).	

$$\mathfrak{r}^{\dagger} = \begin{pmatrix} 973 & 32 & 89 \\ & 1192 & -1 \\ & 583 \end{pmatrix} \times 10^{-4} \, (\text{\AA})^2 \qquad \omega^{\dagger} = \begin{pmatrix} 74 \cdot 0 & -0 \cdot 6 & -9 \cdot 3 \\ & 7 \cdot 0 & -1 \cdot 1 \\ & 59 \cdot 3 \end{pmatrix} \, (^{\circ})^2$$

Direction cosines ( $\times 10^4$ ) for principal axes relative to orthogonal axes:

Eigenvalue c 0·1196 (Å) <sup>2</sup> 0·0989 0·0563	<i>a</i> 	<i>b</i> - 9898 - 1405 243	c* 67 2159 9764		Eigenvalue ω 78·6 (°) <sup>2</sup> 54·7 4·3	8 - 4	954 451 117	b - 2 259 9997	c* - 4452 - 8951 230
			$U_{ij}$	$(Å)^2 \times 10^3$	*†‡				
	C(1)	U <sub>11</sub> 133 -17	U <sub>22</sub> 119 -45	U <sub>33</sub> 103 - 18	$U_{12} = \frac{3}{-3}$	U <sub>13</sub> 15 2	$U_{23} - 1$		
	C(2)	107 	147 43	73 12	-15 -4	10 - 5	7 1		
	C(3)	137 1	239 14	119 17	65 9	12 6	-30 4		
	C(4)	107 - 29	373 - 7	93 7	-42 -11	3 0	$-20\\6$		

† Referred to the directions of the orthogonal axes  $a, b, c^*$ .

<sup>‡</sup> The numbers below the  $U_{ij}$  are the differences (× 10<sup>3</sup>) of  $U_{ij}$  derived from  $B_{ij}$  and those calculated from rigid-body parameters. The r.m.s. difference is 0.017 (Å)<sup>2</sup>.

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Table 6. Comparison of observed and final calculated structure factors The asterisked line gives l, h; subsequent lines give  $k, 25F_o, 25F_c$ .)

• .	•,,,,	311									
	-1939	-1927	12	2728	3133	•	784	103	•	-857	-246
. !	411	374 1296	,	-355	-220	3	-694	-765	-8	473	767
	-132	-100 -1258	٠	-245	-311	4	-379	-342		-728	-496
15	-106 -250	-108 -227	13 5	-130 -742	-20 -702	11 3	-192 -858	262 -852	:	231 1378	46 1371
12	-133	-202	10	170	196	•	245	<b>3</b> 14	é	-463	-398
•	5.59	-89	,	306	392	,	203	#18	3	-456	-832
• 12	•77	-26	10	.96	18		-80	-41	•	4114	-105
• 0	7-106	-214	•	-120	-21		184	140			
• 0	-214	-214				•		•••	•	314	347
• 3	9.170	188	•	276	297	2	59	24	0	97	113
•	10 -87	-02	•	***							
•	-131	•71	5	116	81	3	-83	-54	1	-113	-116
	13 65	25									
• 1	-82	-49									
10	-386	145	8 0	-250	-283 -296	•	-460	-435	4	-53	-99
. 1	103	98	5	190	150	3	-972	-964	1	397	471
10	89 739	122 713	8	255 -1237	244 -1234	6	132	67	•	1117	1125
• •	3275	312	,	338	271	3	-108	-161	1	-113	-148
14	-81	-55 -48 J	10	162	224	6	-63	-51	4	-834	-846
• 1	5379	383	5	-143	-183	3	224	298	1	129	
	<b>9</b> 265	265	2	-79	-74	0	336	292	-		•••
• 1	-244 8	-253	5	-373	-397	3	-175	-171	1	45	42
• 1	-292	-320	4	-188	-538	2	132	108	0	169	183
• •	10 74	-47	3	-56	-67	1	234	218			
• i 1	11 70	69	-	••••		•	<i>.</i>	119	•	"	114
• 1	12 -77	-59									
13	-114	-120	11	-313 -152	-348	•	-360	-346	,	-272	-257
• 12	-2 76		à	-145	-118	•	348	331	•	-221	-260
15	-3 74	95	11	-73	-49	•	294	269	,	220	124
• 1	-4 89	14	8	211	155		-001	-099		-231	-959
• 1		-359	0	320	201						-176
• 1	-6 3:4	2.5		09 318		3	-782	-712	1	310	267
• 1	-7	-67			•••		3/	74	•	220	273
• 1	-, 122	139	3	-232	-258	1	166	147			
. 1	-9 -7		•	-57	-4/	•	-454	-470	4	-297	-324
- 1	-154	532	,	-159	-127	,	-230	-282	3	-78	-66
	-11 -11	-32	2	154	60	٥	287	282			
• 2	111	74	,	-84	-94	i	69	122			
10	1159	1176	6 2	134 2056	199 2015	4	-462	-479	•	-297	-315
11	-64	-026	7	-196	-189	5	-267	-280	3	-103	-91
12	-795	a) •777	10	136	106	a	-151	-116	•	-277	-259
• 2	3200	30 .	,	-151	-66	,	-502	-470	3		28
• 2 11	-24: 4 -54	-292. -75		102	124	-			•		
• 2	- 679	. 18,	ů	-645	-613	•	687	-10	•	-207	-261
• 2	-73	• •	5	-45	-29	3	-266	-243	1	-91	-56
• ;	7 189	225	•	304	204	2	117	120	8	-288	-289
• 2	134	13>	2	169	193	4	183	181			
• • •	<b>.</b> ,	-92									
	17 17	67	2	-02	103	•	-105	-94			
6	-1/0	-31	4	-78	- 57						

.

# Table 6 (cont.)

· 2 11 3	-1 305 -207	360 -293	•	229 4812	275	,	57	48	5	-682	-762
• 2	-2 154 -25	162	•	84	151	•	-354	-991		546	394
• 2	-3 -101 -347	-82	,	-165	-157	5	-309	-814	3	-578	-529
• 2	-4 205	193	6	-432	-391	4	-655	-642	2	212	135
• 29	-5 225	245	,	-285	-253	3	-64	-62	1	654	-634
10	-6 -84 -680	-107 -681	8 0	69 -292	37	6	341	343	4	-334	-289
• 2	-7 -96	-80	1	•79	-67						
• 2	175 -9	172	4	269	282	0	-251	-254			
• 2	-10 -191	-109	2	168	175	5	342	338 196	3	137	112
• 2	-12 -105 -13	-99									
• 3 10	73 0	17		-144	-147						
• 3	-348	-312		- 140	- 245	•			•	-273	-241
• 3	-170	-15/	,	98	82	,	283	250	3	554	556
• 3	-261 3	-243		122	144	6	225	201	•	263	231
11 3 • 3	114 -608 4	107	9 1	70 136	117 113	,	78	122	5	-206	-209
12	75 114	19	10	156 65	204	8	208	235	4	-92	-55
11	87	24	7	130	144	5	128	124	1	114	111
• 3 13	-96 7 -137	-30	6 7	-183 -193	-190 -173	4	-192 -253	-208	1	•0	101
• 3 10 • 3	8-64	-22	0	240	211				-		
13	-132	-13/ -227	11 3	-331 129	-343 141	9	-234	-250	,	-255	-248
12	-111 -223	-125 -252	10 2	-233 -56	-249 -62	8 6	-187 280	-192 323	6	-102	-93
• 3 11 • 3	-3 64 -4	63	,	82	60	5	160	169			
10	30e	45 321	6 0	197 -638	192 -626	6	176	128	4	614	621
• 3	-6 171	784	3	-362	-419	1	125	120			
4 • 3	-381 -7	-365	10 2	-109	-75	8 Ú	144 276	206	6	-179	-147
• 3 10	-5 82	190	7	169	157	3	84 - 79	111	1	95	57
• 3	-9-253	-232	5	- 359	-371	-			•	1/0	
• 3	-127	-115	4	-128	-141	C	136	203			
• •	0 -178	-184	4	-155	-43	2	-103	-86			
* 4	1 -136 2	-10/	,	-247	-259	3	-340	-319	1	227	248
10 2	111 40	127 23	8	81	98	6	-256	-213	4	-228	-297
9	99 -391	33 -350	7	149	133	,	54	37	3	-204	-211
	5.0	52	4	-124	-101	2	-188	-161	0	-69	-36
• 4	111 6 191	200	1	-70 171	-60						
3	7 91	100	1	110	110						
• 4	, 68 9 • 75	31.		- 16	- 4 4						
11	-1 149	143	9	143	197	5	-309	-305	3	111	107
• 4 10	-2 215	236	6	-197	-166		-345	-337	2	815	815
• •	-3 -112	-100	,	-205	-221	3	124	147	1	-42	- 33
• 4	-4 -87 24	58+ 58	6	-140	-149	4	-299	-288	2	-186	-164
	-5 90	56	9	114	109	,	-135	-108	,	-311	-300
• 4	-6132	134	4	-152	-107	2	-251	-221	0	-427	-393
	-7 129 -8	97	3	-334	-293	1	-226	-198			
• •	-9 <sup>52</sup>	37	0	-87 213	-88	3	140	134		-54	- 41
	-10 226	147	4	77	144	2	59	78	-	63	92
• 5	76	67	4	146	142	2	60	74	0	-57	-55
• 5	2 2 -179	94 -180	3	149 -143	133 -131	1 0	-45 85	-64 78			
• •	3 127 4	141	7	80	46	,	-98	-89	1	49	59
• •	5 <sup>61</sup>	69 .01	6	66 -84	68 - 74	4	59 43	34	•	82	74
	-152	-165	•	-18	-126	•	72	54			
• 5		•39	1	79	81						
• ;	•	-24									

-1 -104	-107	,	-47	-29	3	-172	-154	1	-25	-14
-2 -123	-130	5	-237	-124						
-3 -106 72	-121 85	,	-65	-50	,	-70	-58	3	-112	-188
-4 64	70	4	119	106	2	38	45			
-5 05	60	,	95	161	,	186	150	3	202	203
· -6	75	4	-151	-150						
135	161	5	-142	-128	3	-87	-83			
5 -8 3 117	101	6'	69	50						
-9-86	-90									
-10	-151									
2 241	283	6	535	587						
465	518									

Table 6 (cont.)

crepancies are probably due to the inaccuracies of the  $U_{ij}$  values for the atom C(4), caused by the interaction of thermal and positional parameters for this atom and its mirror-related atom C(4'). For anthracene, however, the axes of libration are not simply related to the molecular axes, except that both axes about which the libration is large are nearly perpendicular to the short axis in the plane of the molecule. Their combined effect is to add libration of the molecule in its own plane to the type of motion found for TCNQ.

There are no abnormal sideways intermolecular contacts in the structure. The closest approach,  $d_{C \cdot \cdot N} =$ 3·47 Å, is shown in Fig.2. There are no other intermolecular approaches of less than 3·57 Å. The *b* cell length seems to be determined by the size of the TCNQ molecule. The evidence for this is that the shortest intermolecular contacts, between stacks, involve only the TCNQ molecules. This explains the relatively small thermal librations and vibrations of the TCNQ molecule as compared with those of the anthracene molecule and the fact that anthracene, but not TCNQ, has librational motion in its own plane.

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#### References

ARNDT, U. W. & PHILLIPS, D. C. (1961). Acta Cryst. 14, 807. BLOW, D. M. (1960). Acta Cryst. 13, 168.





- BUSING, W. R. & LEVY, H. A. (1964). Acta Cryst. 17, 142. CRUICKSHANK, D.W. J. (1956). Acta Cryst. 9, 754.
- CRUICKSHANK, D.W. J. & SPARKS, R. A. (1960). Proc. Roy. Soc. A 258, 270.
- HOERNI, J. A. & IBERS, J. A. (1954). Acta Cryst. 7, 744.
- LONG, R. E., SPARKS, R. A. & TRUEBLOOD, K. N. (1965). Acta Cryst. 18, 932.
- MELBY, L. R., HARDER, R. J., HERTLER, W. R., MAHLER, W., BENSON, R. E. & MOCHEL, W. E. (1962). J. Amer. Chem. Soc. 84, 3374.
- SCHOMAKER, V., WASER, J., MARSH, R. E. & BERGMAN, G. (1959). Acta Cryst. 12, 600.
- WILLIAMS, R. M. & WALLWORK, S. C. (1967). Acta Cryst. 22, 899.