# Molecular Complexes Exhibiting Polarization Bonding. XI. The Crystal and Molecular Structure of the 7,7,8,8-Tetracyanoquinodimethane-Anthracene Complex 

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(Received 12 April 1967)


#### Abstract

The $1: 1$ complex formed between $7,7,8,8$-tetracyanoquinodimethane and anthracene crystallizes as deep green needles, space group $C 2 / m$, with two molecules of complex in the unit cell of dimensions $a=$ $11 \cdot 476 \pm 0.011, b=12 \cdot 947 \pm 0.013, c=7 \cdot 004 \pm 0 \cdot 007 \AA, \beta=105 \cdot 4 \pm 0 \cdot 2^{\circ}$. The structure was solved by Patterson methods and refined by least-squares methods using 469 independent reflexions measured on a linear diffractometer. The final residual, $R$, was $0 \cdot 111$. The component molecules are stacked alternately in infinite columns which are parallel to the $c$ axis with the mean plane of the molecules perpendicular to the stacking axis and a mean perpendicular separation of $3 \cdot 50 \AA$. The central six-membered rings of the two types of molecule are orientated at $30^{\circ}$ to each other.


## Introduction

The $1: 1$ complex formed between 7,7,8,8-tetracyanoquinodimethane (TCNQ) and anthracene was first described by Melby, Harder, Hertler, Mahler, Benson \& Mochel (1962). The investigation of the crystal and molecular structure of the complex was carried out to provide comparison with the structures of complexes of aromatic molecules with tetracyanoethylene (TCNE) as exemplified by the naphthalene-TCNE complex described in Part IX (Williams \& Wallwork, 1967).

## Experimental

Deep-green needles of the complex were deposited after hot equimolar solutions of the components in tetrahydrofuran were mixed and allowed to cool very slowly. Microanalyses for carbon, nitrogen and hydrogen confirmed the $1: 1$ ratio of the component molecules in the complex. In plane polarized light the crystals exhibited strong pleochroism. The colours ranged from deep green to a straw colour as the plane of polarization of the light was rotated and maximum absorption occurred when the electric vector of the polarized light was vibrating parallel to the needle axis.

Oscillation and Weissenberg X-ray photographs were taken, using $\mathrm{Cu} K \alpha$ radiation $(\lambda=1.542 \AA)$, with the crystal rotating about the $b$ and the $c$ (needle) axes. The cell dimensions, determined from the Weissenberg photographs, were refined with the use of a linear diffractometer (Arndt \& Phillips, 1961), and intensity data for the $h 0 l, h 1 l, \ldots, h 11 l$ and the $h k 0, h k 1, \ldots, h k 6$ reciprocal lattice levels were then collected with Mo $K \alpha$ radiation. A total of 469 independent reflexions with $\theta<32.5^{\circ}$ were observed out of a possible total of 1800 accessible reflexions. Intensity measurements were

[^0]made with a scintillation counter and a pulse height selector. The moving-crystal stationary-counter technique was employed. The crystal used was approximately cylindrical with a radius of 0.2 mm . No absorption corrections were applied to the data.

## Crystal data

$\mathrm{C}_{8} \mathrm{H}_{4}(\mathrm{CN})_{4} \cdot \mathrm{C}_{14} \mathrm{H}_{10}, \quad M=382 \cdot 43$. Monoclinic, $a=$ $11.476 \pm 0.011, b=12.947 \pm 0.013, c=7.004 \pm 0.007 \AA$, $\beta=105 \cdot 4^{\circ} \pm 0 \cdot 2^{\circ}, U=1003 \AA^{3}, D_{m}=1 \cdot 25 \mathrm{~g} . \mathrm{cm}^{-3}, Z=$ $2, D_{c}=1.26 \mathrm{~g}_{\mathrm{cm}}{ }^{-3}, F(000)=396$. Мо $K \alpha(\lambda=0.7107 \AA)$, $\mu=0.84 \mathrm{~cm}^{-1}$. Absent spectra, $h k l$ when $h+k$ odd. Space group C2 (no.5), Cm (no.8) or $C 2 / m$ (no.12). Negative piezoelectric and pyroelectric tests indicate $C 2 / m$ and this is confirmed by refinement of the structure.

## Determination and refinement of the structure

A trial structure was deduced from the $h k 0$ and $h 0 l$ Patterson projections. The structure was partially refined through structure-factor, electron-density projection calculations to the stage where the reliability index,

$$
R=\Sigma \| F_{o}\left|-\left|F_{c}\right|\right| / \Sigma\left|F_{o}\right|,
$$

was 0.21 and 0.16 for the $h k 0$ and $h 0 l$ zones respectively. Structure factors were then calculated for the $h k l$ reflexions followed by a three-dimensional Fourier elec-tron-density summation, from which improved atomic coordinates were obtained.

Refinement was continued by the block-diagonal least-squares method with a program written by Dr R.D.Diamand. The weighting scheme employed was $\forall w=1$ if $\left|F_{0}\right|<F^{*}$ otherwise $\backslash w=F^{*}| | F_{0} \mid$, and the scattering factors of Hoerni \& Ibers (1954) for carbon and nitrogen were used. The reflexions too weak to be observed were omitted from the analysis. Hydrogen atom parameters were not included. Four cycles of structure-factor least-squares using individual isotropic temperature factors improved $R$ to $0 \cdot 165$. Refinement
was then continued allowing the temperature factors of the atoms to vary individually and anisotropically. After four cycles of least-squares the refinement converged with $R=0 \cdot 111$ for observed reflexions. The comparison of observed and final calculated structure factors is shown in Table 6.

The final atomic coordinates and their estimated standard deviations (derived from the least-squares normal equations block-diagonal matrix) are summarized in Table 1. The anisotropic temperature factors are given in Table 2.

Table 1. Final positional parameters (upper line) and their estimated standard deviations (lower line)

| C(1) | $10^{4} x / a$ | $10^{4} y / b$ | 104z/c |
| :---: | :---: | :---: | :---: |
|  | 0000 | 1088 | 0000 |
| C(2) | 1077 | 0551 | 0475 |
|  | 7 | , | 12 |
| C(3) | 2214 | 1127 | 0938 |
|  | 10 | 12 | 16 |
| $\mathrm{C}(4)$ | 3228 | 0499 | 1382 |
| C(5) | 1271 | 13 0000 | 5509 |
|  | 9 |  | 14 |
| C(6) | 0617 | 0975 | 5241 |
|  | 6 | 5 | 11 |
| C (7) | 2506 | 0000 | 6026 |
| C (8) | 3198 | 0929 | 6325 |
|  |  | 6 | 12 |
| N | 3764 | 1657 | 6574 |
|  | 6 | 6 | 12 |

Table 2. Final atomic thermal parameters ( $\times 10^{4}$ )
$B_{i j}$ are coefficients in the temperature factor expression $\exp \left[-\left(B_{11} h^{2}+B_{22} k^{2}+B_{33} l^{2}+B_{12} h k+B_{13} h l+B_{23} k l\right)\right]$.

|  | $B_{11}$ | $B_{22}$ | $B_{33}$ | $B_{12}$ | $B_{13}$ | $B_{23}$ |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: |
| $\mathrm{C}(1)$ | 198 | 87 | 367 | 000 | 205 | 000 |
| $\mathrm{C}(2)$ | 145 | 122 | 367 | -55 | 140 | -37 |
| $\mathrm{C}(3)$ | 230 | 297 | 443 | -214 | 235 | -116 |
| $\mathrm{C}(4)$ | 217 | 432 | 433 | -152 | 156 | -65 |
| $\mathrm{C}(5)$ | 116 | 69 | 268 | 000 | 129 | 000 |
| $\mathrm{C}(6)$ | 132 | 71 | 306 | 13 | 124 | 9 |
| $\mathrm{C}(7)$ | 122 | 72 | 318 | 000 | 83 | 000 |
| $\mathrm{C}(8)$ | 114 | 91 | 442 | -1 | 123 | 28 |
| N | 149 | 100 | 634 | -54 | 120 | -3 |

Sections of the three-dimensional electron-density distribution and of the difference electron-density distribution were then calculated. These syntheses showed no unexpected features, confirmed the marked anisotropy of the anthracene molecule and clearly indicated the positions of the hydrogen atoms of the TCNQ molecule.

Least-squares planes through each of the molecules were calculated by the modification (Blow, 1960) of the method of Schomaker, Waser, Marsh \& Bergman (1959). All of the atoms, other than hydrogen, in each molecule were used to define the planes and each atom used was given unit weight. The pertinent parameters of the planes, and the deviations of the atoms from them, are summarized in Table 3.

Table 3. Molecular least-squares planes
TCNQ
Anthracene
Perpendicular distance in $\AA$ from the origin
3.502

Direction cosines of normal relative to orthogonal axes

| $a$ | 0.2533 | 0.2612 |
| :--- | ---: | ---: |
| $b$ | 0.0000 | 0.0000 |
| $c^{*}$ | -0.9674 | -0.9653 |

Perpendicular distance in $\AA$ of atoms from the mean planes

| $\mathrm{C}(5)$ | 0.013 | $\mathrm{C}(1)$ | 0.000 |
| :--- | ---: | ---: | ---: |
| $\mathrm{C}(6)$ | 0.010 | $\mathrm{C}(2)$ | -0.010 |
| $\mathrm{C}(7)$ | 0.009 | $\mathrm{C}(3)$ | 0.006 |
| $\mathrm{C}(8)$ | 0.001 | $\mathrm{C}(4)$ | -0.001 |
| $\mathrm{~N}(1)$ | -0.008 |  |  |

The anisotropic thermal parameters of the atoms were used in an analysis of the rigid-body translational and librational motion (Cruickshank, 1956) of each molecule. The results of the analyses are given in Table 4 for the TCNQ molecule and in Table 5 for the anthracene molecule. The r.m.s. discrepancies between the observed and calculated $U_{i j}$ values are 0.004 and $0.017 \AA^{2}$ for the TCNQ and anthracene molecules respectively. The method of Busing \& Levy (1964) was used to correct the bond lengths of the TCNQ molecule for libration. The corrected bond lengths are illustrated in Fig. $1(a)$ together with the uncorrected bond lengths and angles and their e.s.d.'s. The bond lengths and angles and their e.s.d.'s for the anthracene molecule are illustrated in Fig.1(b). No attempt was made to correct these for librational motion because it was so extreme as to cause atoms $C(4)$ and $C\left(4^{\prime}\right)$ to be unresolved in one electron-density peak.

## Description and discussion of the structure

The structure consists of infinite columns of alternate TCNQ and anthracene molecules stacked in a plane-to-plane manner in columns parallel to the $c$ axis of the unit cell with the mean molecular planes perpendicular to this axis. Each cell contains two columns of molecules and the structure, projected along the $c$ axis, is illustrated in Fig.2. The mode of molecular overlap is such that the long axes of each molecule are parallel, and their centres are superimposed. The central hexagonal ring systems of the two types of molecule are orientated at $30^{\circ}$ to each other, thereby avoiding direct overlap of the majority of the atoms. This relative orientation is not that for the optimum overlap of the $\pi$-orbital systems of the molecules but the structure is such that the aromatic system of the TCNQ molecule is overlapped by the entire aromatic system of the anthracene molecule. There is no suggestion of specific interaction involving $\mathrm{C}=\mathrm{C}$ bonds as found in the structure of the naphthalene-TCNE complex (Williams \& Wallwork, 1967).

The most surprising aspect of the structure is the relatively large intermolecular spacing ( $3 \cdot 50 \AA$ ) of the


Fig. 1. Bond lengths ( $\AA$ ) and bond angles, with (in parentheses) their standard deviations, for (a) TCNQ, (b) anthracene. Values for TCNQ in square brackets are bond lengths corrected for libration.

Table 4. Rigid-body thermal parameters for the TCNQ molecule
(For the molecule centred at $0,0, \frac{1}{2}$ ).

$$
\tau \dagger=\left(\begin{array}{ccc}
746 & -1 & 729 \\
& 547 & 2 \\
& & 536
\end{array}\right) \times 10^{-4} \AA^{2} \quad \omega \dagger=\left(\begin{array}{rrr}
32 \cdot 5 & 0.1 & 8 \cdot 1 \\
& 9 \cdot 3 & -0.7 \\
& & 8 \cdot 1
\end{array}\right)\left({ }^{\circ}\right)^{2}
$$

Direction cosines $\left(\times 10^{4}\right)$ for principal axes relative to orthogonal axes:

| Eigenvalue | $a$ | $b$ | $c$ |  | Eigenvalue |  | $a$ |  | $b$ | $c^{*}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\tau 0.0769 \AA^{2}$ | 9543 | -7 |  |  | (b 34.8 ${ }^{\circ}{ }^{\circ}{ }^{2}$ |  | 9596 |  | -32 | 2814 |
| 0.0547 | 139 | -9988 | - |  | $9 \cdot 6$ |  | -450 |  | -9888 | 1425 |
| 0.0513 | -2984 | -489 | 9593 |  | $5 \cdot 8$ | -2778 |  |  | 1494 | 9490 |
| $U_{i j}(\AA)^{2} \times 10^{3} \dagger \ddagger$ |  |  |  |  |  |  |  |  |  |  |
|  |  | $U_{11}$ | $U_{22}$ | $U_{33}$ | $U_{12}$ | $U_{13}$ |  | $U_{23}$ |  |  |
|  | C(5) | 79 | 56 | 71 | -2 | 3 |  | 0 |  |  |
|  |  | 1 | 4 | 0 | 6 | 2 |  | 2 |  |  |
|  | C(6) | 75 | 59 | 59 | 0 | 6 |  | 0 |  |  |
|  |  | $-7$ | 1 | 3 | 0 | 2 |  | 0 |  |  |
|  | C(7) | 76 | 68 | 75 | 0 | 2 |  | 1 |  |  |
|  |  | 2 | $-7$ | -2 | 0 | -6 |  | -1 |  |  |
|  | C(8) | 80 | 77 | 102 | -8 | -5 |  | 2 |  |  |
|  |  | $-9$ | 0 | 0 | 6 | 1 |  | 9 |  |  |
|  | N(1) | 89 | 86 | 147 | -16 | -16 |  | 2 |  |  |
|  |  | 8 | -1 | -1 | -4 | 0 |  | $-3$ |  |  |

$\dagger$ Referred to the directions of the orthogonal axes $a, b, c^{*}$.
$\ddagger$ The numbers below the $U_{i j}$ are the differences $\left(\times 10^{3}\right)$ of $U_{i j}$ derived from $B_{i j}$ and those calculated from rigid-body parameters. The r.m.s. difference is $0.004(\AA)^{2}$.
molecules within one stack, presumably due to the unfavourable relative orientation of the components. The four shortest approaches between atoms of adjacent molecules of the same column $(\mathrm{C}(1)-\mathrm{C}(6) 3 \cdot 57, \mathrm{C}(2)-$ $\mathrm{C}(6) 3.58, \mathrm{C}(2)-\mathrm{C}(5) 3.58, \mathrm{C}(4)-\mathrm{C}(8) 3.58 \AA)$ are, if anything, greater than the expected van der Waals distances. This implies that any charge-transfer interaction between the molecules is weak. This does not conflict with the existence of a low-energy, charge-transfer band in the spectrum of the complex, which causes its deep green colour. The fairly low ionization potential


Fig. 2. The molecular arrangement, viewed along the $c$ axis.
of anthracene and the high electron affinity of TCNQ inevitably produce an excited state of fairly low energy, but this will not stabilize the ground state of the complex by admixture of the two wave functions if the symmetries are inappropriate for overlap.

Both molecules are planar within experimental error and the molecular planes are inclined to each other at the small angle of $\frac{1}{2}^{\circ}$. The structure projected along the $b$ axis is illustrated in Fig. 3.

The libration-corrected dimensions of the TCNQ molecule are in good agreement with the dimensions of the uncomplexed molecule (Long, Sparks \& Trueblood, 1965). The r.m.s. discrepancy between the two determinations is $0.007 \AA$, much less than one e.s.d. No bond length differs by more than $1 \frac{1}{2}$ e.s.d. The maximum difference in the bond angles is also $1 \frac{1}{2}$ e.s.d., the r.m.s. difference being one degree.

The dimensions of the anthracene molecule [apart from the bonds and angles involving the atom $\mathrm{C}(4)$ ] are in reasonable agreement with the weighted mean of the bond lengths determined by various authors and discussed by Cruickshank \& Sparks (1960). The dimensions involving the atom $C(4)$ are invalid for the reasons discussed above.

The thermal parameters, $U_{i j}$ of the TCNQ molecule appear to be adequately explained on the basis of rigidbody librations and vibrations. The principal axes of both libration and of vibration correspond to within a few degrees of the natural molecular axes and there is appreciable libration only about the long axis in the plane of the molecule. The thermal parameters, $U_{i j}$, of the anthracene molecule are reasonably well explained on the basis of rigid-body motions. The dis-

Table 5. Rigid-body parameters for the anthracene molecule
(For the molecule centred at $0,0,0$ ).

$$
\tau \dagger=\left(\begin{array}{rrr}
973 & 32 & 89 \\
& 1192 & -1 \\
& & 583
\end{array}\right) \times 10^{-4}(\AA)^{2} \quad \omega \dagger=\left(\begin{array}{rrr}
74 \cdot 0 & -0 \cdot 6 & -9 \cdot 3 \\
& 7 \cdot 0 & -1 \cdot 1 \\
& & 59 \cdot 3
\end{array}\right)\left({ }^{\circ}\right)^{2}
$$

| Eigenvalue | $a$ | $b$ | $c^{*}$ |  | Eigenvalue |  | $a$ | $b$ | $c^{*}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| c $0 \cdot 1196(\AA)^{2}$ | -1425 | -9898 | -67 |  | ف $\left.78.6{ }^{( }\right)^{2}$ |  | 8954 | -2 | -4452 |
| 0.0989 | 9662 | -1405 | 2159 |  | 54.7 |  | -4451 | 259 | -8951 |
| 0.0563 | -2147 | 243 | 9764 |  | $4 \cdot 3$ |  | 117 | 9997 | 230 |
| $U_{i j}(\AA)^{2} \times 10^{3} \dagger \ddagger$ |  |  |  |  |  |  |  |  |  |
|  |  | $U_{11}$ | $U_{22}$ | $U_{33}$ | $U_{12}$ | $U_{13}$ | $U_{23}$ |  |  |
|  | C(1) | 133 | 119 | 103 | 3 | 15 | $-1$ |  |  |
|  |  | -17 | -45 | -18 | -3 | 2 |  |  |  |
|  | C(2) | 107 | 147 | 73 | -15 | 10 | -7 |  |  |
|  |  | -15 | -43 | 12 | -4 | -5 | -1 |  |  |
|  | C(3) | 137 | 239 | 119 | -65 | 12 | -304 |  |  |
|  |  | -1 | 14 | 17 | -9 | 6 |  |  |  |
|  | C(4) | 107 | 373 | 93 | -42 | 3 | -206 |  |  |
|  |  | -29 | -7 | 7 | -11 | 0 |  |  |  |

$\dagger$ Referred to the directions of the orthogonal axes $a, b, c^{*}$.
$\ddagger$ The numbers below the $U_{i j}$ are the differences $\left(\times 10^{3}\right)$ of $U_{i j}$ derived from $B_{i j}$ and those calculated from rigid-body parameters. The r.m.s. difference is $0.017(\AA)^{2}$.

Table 6. Comparison of observed and final calculated structure factors
The asterisked line gives $l, h$; subsequent lines give $k, 25 F_{o}, 25 F_{c}$.)


Table 6 (cont.)


Table 6 (cont.)

crepancies are probably due to the inaccuracies of the $U_{i j}$ values for the atom C(4), caused by the interaction of thermal and positional parameters for this atom and its mirror-related atom $C\left(4^{\prime}\right)$. For anthracene, however, the axes of libration are not simply related to the molecular axes, except that both axes about which the libration is large are nearly perpendicular to the short axis in the plane of the molecule. Their combined effect is to add libration of the molecule in its own plane to the type of motion found for TCNQ.

There are no abnormal sideways intermolecular contacts in the structure. The closest approach, $d_{\mathrm{C} \cdot \cdots \mathrm{N}}=$ $3.47 \AA$, is shown in Fig. 2. There are no other intermolecular approaches of less than $3 \cdot 57 \AA$. The $b$ cell length seems to be determined by the size of the TCNQ molecule. The evidence for this is that the shortest intermolecular contacts, between stacks, involve only the TCNQ molecules. This explains the relatively small thermal librations and vibrations of the TCNQ molecule as compared with those of the anthracene molecule and the fact that anthracene, but not TCNQ, has librational motion in its own plane.

We are indebted to the Director of the Manchester University Department of Computer Science for allowing computing facilities on the Atlas computer, to the Director of the Oxford University Computing Laboratory for allowing computing facilities on the Mercury computer, and to M.M.Harding, R.D.Diamand, J.S. Rollett and J.H. Rayner for the use of their computer programs. We also thank Dr J.S. Rollett for assistance with the Mercury computing and the Science Research Council for a maintenance grant (to R.M.W.) and for providing the linear diffractometer.

## References

Arndt, U. W. \& Phllips, D. C. (1961). Acta Cryst. 14, 807. Blow, D. M. (1960). Acta Cryst. 13, 168.


Fig.3. The molecular arrangement, viewed along the $b$ axis. Atoms belonging to molecules centred at $y=0$ and $y=0.5$ are represented by full and open circles respectively.

Busing, W. R. \& Levy, H. A. (1964). Acta Cryst. 17, 142. Cruickshank, D. W. J. (1956). Acta Cryst. 9, 754.
Cruickshank, D. W. J. \& Sparks, R. A. (1960). Proc. Roy. Soc. A 258, 270.
Hoern, J. A. \& Ibers, J. A. (1954). Acta Cryst. 7, 744.
Long, R. E., Sparks, R. A. \& Trueblood, K. N. (1965). Acta Cryst. 18, 932.
Melby, L. R., Harder, R. J., Hertler, W. R., Mahler, W., Benson, R. E. \& Mochel, W. E. (1962). J. Amer. Chem. Soc. 84, 3374.
Schomaker, V., Waser, J., Marsh, R. E. \& Bergman, G. (1959). Acta Cryst. 12, 600.

Williams, R. M. \& Wallwork, S. C. (1967). Acta Cryst. 22, 899.


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